

Calculating eigenvalues of non-hermitian
matrices to determine resonance energies and
widths

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and

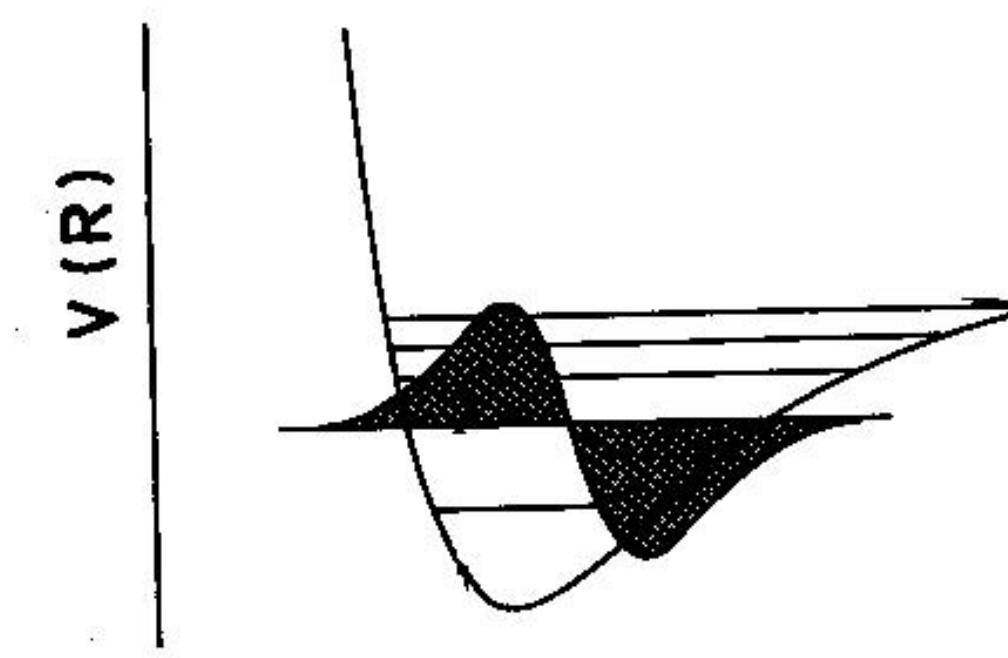
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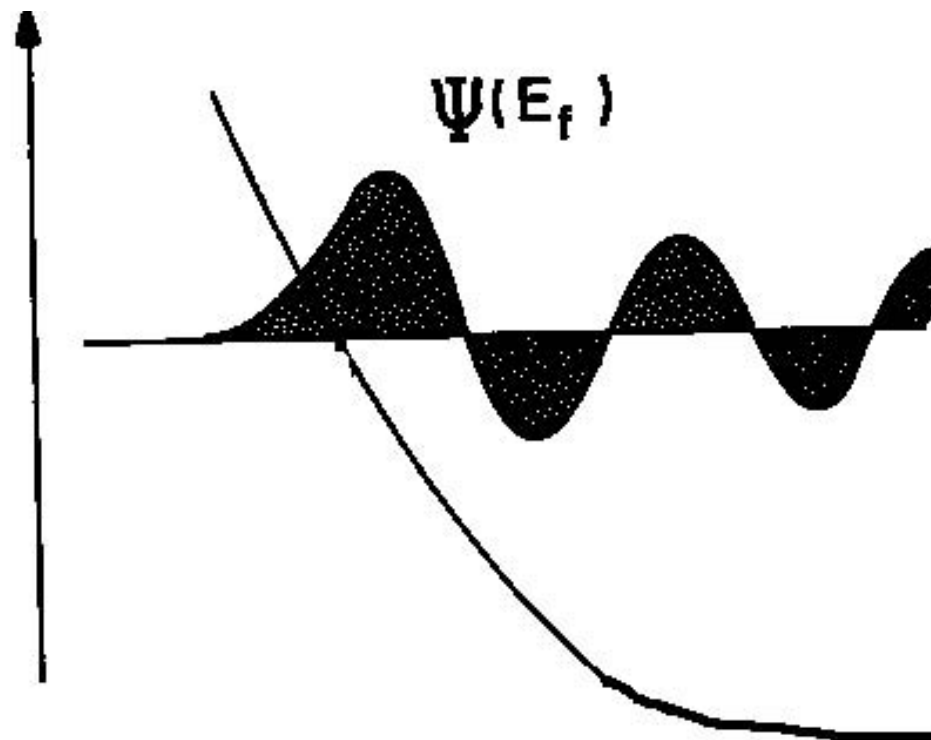
What is a resonance?

Solutions of the Schroedinger equation

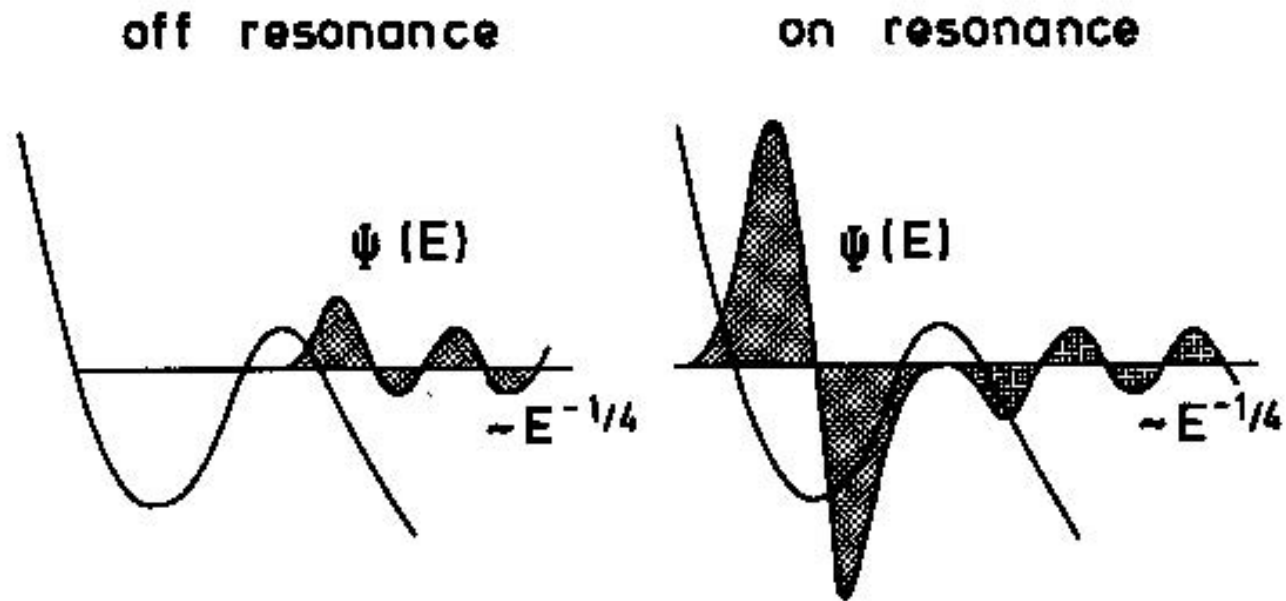
A bound state



A scattering state



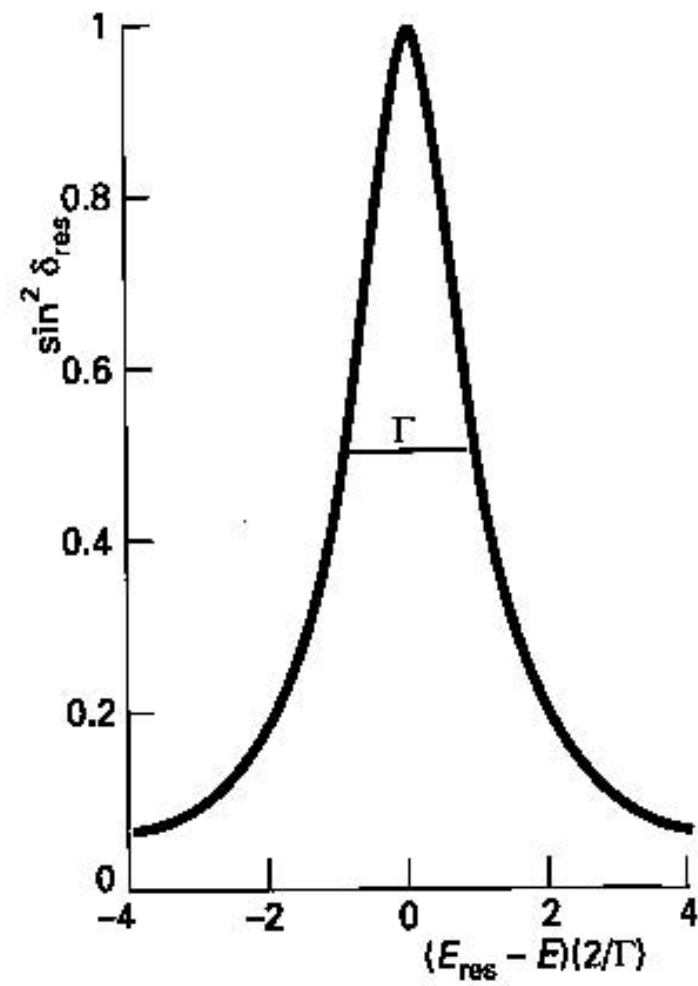
A resonance state



A molecule or complex with a finite lifetime

Resonances are the continuation of the bound states into the continuum

A resonance peak in a spectrum looks like:



It is possible to show that if outgoing boundary conditions are imposed

i.e.

$$\psi \sim e^{+ikr}$$

$$k = \sqrt{2mE}$$

there are solutions of the Schroedinger equation with complex energies

$$E = E_r - i\frac{\Gamma}{2}$$

energy at which resonance occurs width of the resonance

with $\Gamma > 0$

Our goal is to find these states

Problem: ψ diverges

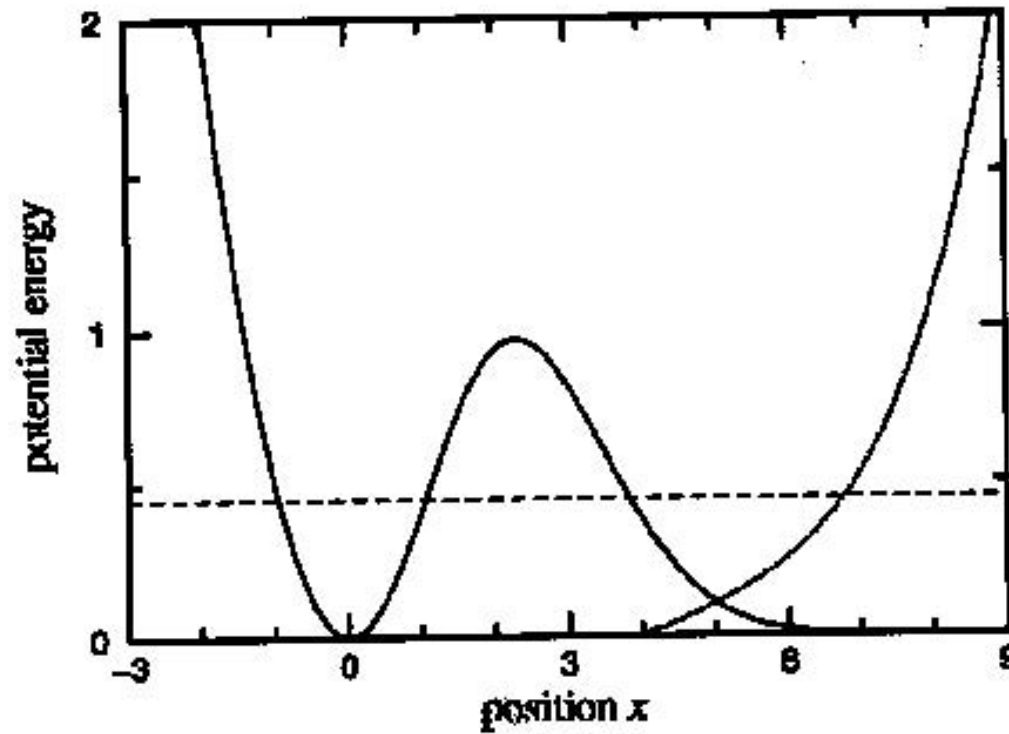
$$E = E_r - i\frac{\Gamma}{2}$$

$$k = k_r - ik_i \quad ; \quad k_i > 0$$

$$e^{+ikr} \rightarrow e^{+irk_r} e^{+rk_i}$$

Standard Approach

$$H \rightarrow H - i\eta W \quad \eta > 0.$$



W kills of the divergence

Choose W so that it does not perturb ψ in the interaction region

η is a parameter that is varied to minimise the perturbation

We discretize to obtain a matrix eigenvalue problem

$$(\mathbf{H} - \mathbf{i}\eta\mathbf{W})\mathbf{Y} = \mathbf{Y}\mathbf{E}$$

How do we compute the (complex) eigenvalues?

Complex symmetric Lanczos with no re-orthogonalization

Cullum and Willoughby

This works quite well

Deficiencies of usual implementation

- Complex matrix-vector products
- No eigenvectors \rightarrow no wavefunctions

New approach

$$(\mathbf{H} - i\eta \mathbf{W}) \mathbf{x}_k = E_k \mathbf{x}_k$$

$$-i\eta \mathbf{W} = \frac{u_k}{2} (\mathbf{I} - \mathbf{D}_1) + \frac{1}{2u_k} (\mathbf{I} - \mathbf{D}_2)$$

where u_k is defined by $E_k = \left(\frac{u_k}{2} + \frac{1}{2u_k} \right)$

Neumaier and Mandelshtam,
Phy. Rev. Lett. 86 5031 (2001)

\mathbf{H} is scaled so that its eigenvalues are between -1 and +1



Quadratic eigenvalue problem

$$(u_k^2 \mathbf{D}_1 - 2\mathbf{H}u_k + \mathbf{D}_2)\mathbf{x}_k = 0 \ .$$

\Downarrow

Linearize

$$\mathbf{U}\hat{\mathbf{x}}_k = u_k\hat{\mathbf{x}}_k$$

where $\hat{\mathbf{x}}_k = \begin{pmatrix} \mathbf{x}_k \\ \mathbf{y}_k \end{pmatrix}$

Advantages of \mathbf{U}

- \mathbf{U} = real
- cost of a matrix-vector product with $\mathbf{U} \sim$ cost of a matrix-vector product with \mathbf{H}
- Ideal for Lanczos because desired eigenvalues are extremal

“Symmetric Indefinite Lanczos Method”

Bai, Ericsson, Kowalski,

Templates for the Solution of Algebraic Eigenvalue Problems

$$\mathbf{U}\hat{\mathbf{x}}_k = u_k\hat{\mathbf{x}}_k \rightarrow \mathbf{UX} = \mathbf{Xu}$$

$$\mathbf{DUX} = \mathbf{DXu}$$

where

$$\mathbf{D} = \begin{pmatrix} \mathbf{D}_2 & \mathbf{0} \\ \mathbf{0} & -\mathbf{D}_1 \end{pmatrix} \text{ is symmetric}$$

Also

$$(\mathbf{DU})^T = \mathbf{U}^T \mathbf{D}^T = \mathbf{U}^T \mathbf{D} = \mathbf{DU}$$

so (\mathbf{DU}) is symmetric

But BOTH (\mathbf{DU}) and \mathbf{D} are not positive definite

The key idea is to use a \mathbf{D} inner product

Basis vectors are orthogonal wrt an indefinite inner product

- loss of linear independence
- failure of the method

It works best for small matrices and for isolated eigenvalues (shift and invert)

For my purpose it fails

Go back to attempting to solve

$$\mathbf{U}\hat{\mathbf{x}}_k = u_k\hat{\mathbf{x}}_k \rightarrow \mathbf{U}\mathbf{X} = \mathbf{X}\mathbf{u}$$

Non-hermitian Lanczos

Two sets of bi-orthogonal vectors formed from the Krylov spaces

$$[\mathbf{q}, \mathbf{U}\mathbf{q}, \mathbf{U}^2\mathbf{q}, \dots] \quad [\mathbf{p}, \mathbf{U}^T\mathbf{p}, (\mathbf{U}^T)^2\mathbf{p}, \dots]$$

The columns of \mathbf{V} and \mathbf{W} are Lanczos vectors

They are generated with three-term recurrence relations

The coefficients are chosen according to the recipe of Freund and Nachtigal

There are different ways to choose the coefficients consistent with the requirement $\mathbf{v}_{i+1}^T \mathbf{w}_{i+1} = 1$

In matrix form the recurrence relations are

$$\mathbf{U}\mathbf{V}_M = \mathbf{V}_{M+1}\mathbf{T}_{M+1}$$

$$\mathbf{U}^T\mathbf{W}_M = \mathbf{W}_{M+1}(\mathbf{\Gamma}^{-1}\mathbf{T}_{M+1}\mathbf{\Gamma})$$

$\mathbf{\Gamma}$ = diagonal matrix related to the way the vectors are normalized

eigenvalues of $\mathbf{T}_M \rightarrow$ eigenvalues of \mathbf{U}

To increase stability and avoid the need to balance we compute eigenvalues of $\mathbf{G}^{-1}\mathbf{T}_M\mathbf{G}$, where \mathbf{G} = diagonal and $\mathbf{G}^{-1}\mathbf{T}_M\mathbf{G}$ is symmetric

Chemist's dream

No reorthogonalization

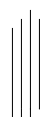
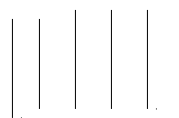
Copied eigenvalues are accurate

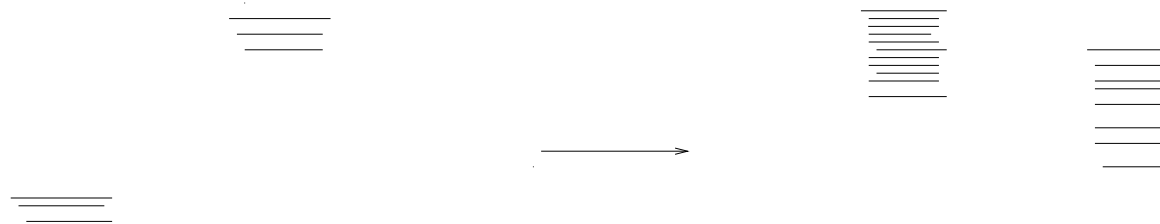
This does not work

Near copies are produced but

norm [copied e'val of largest norm - copied e'val of smallest norm]

increases as we increase the number of matrix-vector products





What is the best estimate of the eigenvalue?

If we use the coupled-two-term approach the situation is somewhat better

Generate \mathbf{p} vectors with \mathbf{v} vectors

Generate \mathbf{q} vectors with \mathbf{w} vectors

The four recursion relations for the four sets of factors each involve only 2 vectors

It is not necessary to do matrix-vector products with both \mathbf{U} and \mathbf{U}^T

Because $\mathbf{U}^T \mathbf{D} = \mathbf{D} \mathbf{U}$

$$\tilde{\mathbf{w}}_{k+1} = \theta_k \mathbf{D} \tilde{\mathbf{v}}_{k+1}$$

$$\mathbf{q}_k = \theta_k \mathbf{D} \mathbf{p}_k$$

θ_k = constant related to the norms of the vectors

Coupled-two-term is better but not good enough (copies are not nearly perfect)

How to compute accurate eigenvalues

- For e'values with norms close to one we identify groups of eigenvalues with almost identical norms

- Average of eigenvalues in a group is \tilde{u}

\tilde{u} = approximate eigenvalue

- Compute approximate right eigenvectors of \mathbf{T}_M by inverse iteration using LU decomposition of $(\tilde{u}\mathbf{I} - \mathbf{T}_M) \rightarrow \mathbf{R}$

- $\mathbf{X}_{\text{Lan}}^{\text{R}} = \mathbf{V}_{\text{ID}} \mathbf{R}$, where ID = optimal iteration depth (between the number of iterations at which first copy occurs and number at which second copy occurs)

- $\mathbf{X}_{\text{Lan}}^{\text{L}} = \mathbf{D} \mathbf{X}_{\text{Lan}}^{\text{R}}$

-

$$\mathbf{U} \mathbf{X} = \mathbf{X} \mathbf{u}$$

$$\mathbf{U} \mathbf{X}_{\text{Lan}}^{\text{R}} \mathbf{Z} = \mathbf{X}_{\text{Lan}}^{\text{R}} \mathbf{Z} \mathbf{u}$$

$$[(\mathbf{X}_{\text{Lan}}^{\text{L}})^{\text{T}} \mathbf{U} \mathbf{X}_{\text{Lan}}^{\text{R}}] \mathbf{Z} = [(\mathbf{X}_{\text{Lan}}^{\text{L}})^{\text{T}} \mathbf{X}_{\text{Lan}}^{\text{R}}] \mathbf{Z} \mathbf{u}$$

$$[(\mathbf{X}_{\text{Lan}}^{\text{L}})^{\text{T}} \mathbf{X}_{\text{Lan}}^{\text{R}}] \text{ is nearly diagonal}$$

- Eigenvalues are computed in groups because to compute them all at once would require storing large $\mathbf{X}_{\text{Lan}}^{\mathbf{R}}$ matrices

Key idea: We use non-hermitian Lanczos ideas to determine a very good subspace and then extract eigenvalues from it

Disadvantage

- We must store Lanczos vectors on disk

Advantage

- No troubles with instability and inaccuracy
- no complex arithmetic

This work has been supported by the
Natural Sciences and Engineering
Research Council of Canada