

Maxwell's equations

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho$$

Coulomb's law

$$\vec{\nabla} \cdot \vec{B} = 0$$

absence of magnetic monopoles

$$\vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$$

Faraday's law of induction

$$\vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} = \frac{4\pi}{c} \vec{J}$$

Ampère's law + Maxwell displacement current

Units - Gaussian

$$[q] \text{ statcoulombs} \quad [\rho] = \text{statcoulombs/cm}^3$$

$$[E] \text{ statvolts/cm} = \text{statcoulombs/cm}^3$$

$$[B] = \text{gauss} \quad [E]/[B] = 1$$

in SI units

$$[E]/[B] = \text{m/s}$$

$$[J] = \text{statamps/cm}^2$$

$$= \text{statcoulombs/cm}^2/\text{s}$$

$$[J]/[\rho] = \text{cm/s}$$

$$qE = \text{force}$$

$$1 \text{ dyne} = \text{statcoulomb}^2/\text{cm}^2$$

$$\frac{1}{8\pi} (E^2 + B^2) = \text{energy density}$$

$$1 \text{ erg/cm}^3 = (\text{statcoulomb/cm}^2)^2$$

$$= \text{dyne/cm}^2$$

$$= \text{gauss}^2$$

Ohm's law $\vec{J} = \sigma \vec{E}$ $[\sigma] = \text{s}^{-1}$

Lorentz force law

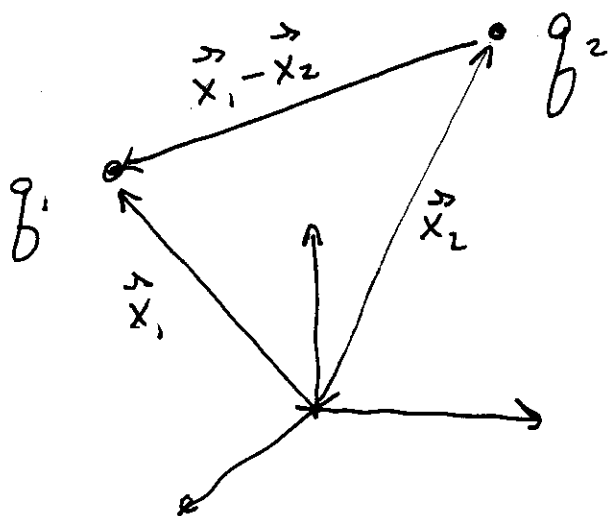
$$\frac{d\vec{p}}{dt} = q \left(\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right)$$

Electrostatics \vec{E}, ρ time independent
 $\vec{J}, \vec{B} = 0$

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$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad \vec{\nabla} \times \vec{E} = 0$$

Begin with Coulomb force law



$$\vec{F}_{12} \text{ on 1 due to 2} = q_1 q_2 \frac{\vec{x}_1 - \vec{x}_2}{|\vec{x}_1 - \vec{x}_2|^3}$$

Electric field

$$\text{field at } \vec{x} \text{ due to 2} = q_2 \frac{\vec{x} - \vec{x}_2}{|\vec{x} - \vec{x}_2|^3} = \vec{E}(\vec{x})$$

$$\vec{F}_{12} \text{ on 1 due to 2} = q_1 \vec{E}(\vec{x}_1)$$

for system of charges

$$\vec{E}(\vec{x}) = \sum_i q_i \frac{\vec{x} - \vec{x}_i}{|\vec{x} - \vec{x}_i|^3}$$

$$q_i \rightarrow \rho(\vec{x}') d^3x'$$

$$\vec{E}(\vec{x}) = \int \rho(\vec{x}') \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} d^3x'$$

formally $\rho(\vec{x}) = q \delta(\vec{x} - \vec{x}_i)$ for pt charge
at $\vec{x} = \vec{x}_i$

We see that

$$\frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} = -\vec{\nabla} \frac{1}{|\vec{x} - \vec{x}'|}$$

Proof $|\vec{x} - \vec{x}'| = (x^2 + x'^2 - 2\vec{x} \cdot \vec{x}')^{1/2}$

$$\vec{\nabla}_x |\vec{x} - \vec{x}'| = \frac{1}{2} ()^{-1/2} \cdot (2\vec{x} - 2\vec{x}')$$

$$\text{So } \nabla_x |\vec{x} - \vec{x}'| = \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3}$$

$$\vec{\nabla} \frac{1}{|\vec{x} - \vec{x}'|} = -\frac{1}{2} ()^{-3/2} (\vec{x} - \vec{x}') = -\frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} \quad \checkmark$$

$$\text{So } \vec{E} = -\vec{\nabla} \Phi \quad \Phi = \int \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$\text{Note } \vec{\nabla} \times \vec{E} = \vec{\nabla} \times \vec{\nabla} \Phi = 0$$

Alternative

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad \vec{\nabla} \times \vec{E} = 0$$

second equation implies $\vec{E} = -\vec{\nabla} \Phi$

$$\text{then } \nabla^2 \Phi = -4\pi\rho$$

Solve using Green's functions

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first solve for δ -fn source

$$\nabla^2 G(\vec{x}, \vec{x}') = -4\pi \delta^3(\vec{x} - \vec{x}')$$

then
$$\Phi(\vec{x}) = \int d^3x' G(\vec{x}, \vec{x}') \rho(\vec{x}')$$

we now show

$$G(\vec{x}, \vec{x}') = \frac{1}{|\vec{x} - \vec{x}'|}$$

i.e.
$$\nabla^2 \frac{1}{|\vec{x} - \vec{x}'|} = -4\pi \delta(\vec{x} - \vec{x}')$$

let $\vec{r} = \vec{x} - \vec{x}'$

$$\nabla^2 \frac{1}{r} = \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \left(\frac{1}{r} \right) = 0 \quad \text{except where } r=0$$

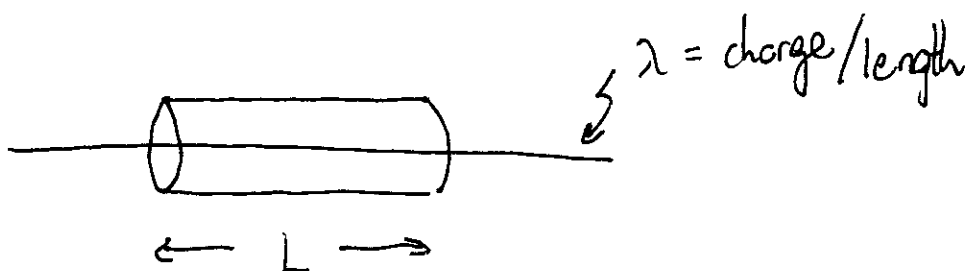
Integrate $\nabla^2 \frac{1}{r}$ over ball of radius $r = R$
centered on origin

$$\begin{aligned} \int \nabla^2 \frac{1}{r} d^3r &= \int (\vec{\nabla} \cdot \vec{\nabla} \frac{1}{r}) d^3r \\ &= \oint \vec{\nabla} \frac{1}{r} \cdot \hat{n} da = \oint \left(-\frac{1}{R^2}\right) R^2 d\Omega \\ &= -4\pi \end{aligned}$$

Gauss's law

$$\int_V \vec{\nabla} \cdot \vec{E} d^3x = \oint_S \vec{E} \cdot \hat{n} da = \int_V d^3x 4\pi\rho$$

e.g. infinite charged
wire

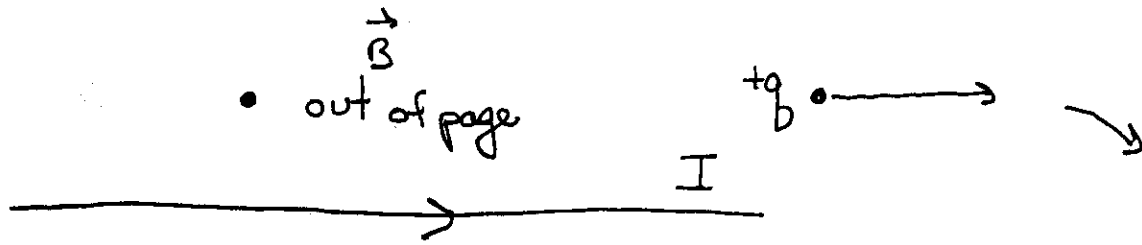


$$\vec{\nabla} = E_s \hat{s}$$

$$2\pi s E_s \cancel{L} = 4\pi \lambda \cancel{L}$$

$$E_s = \frac{2\lambda}{s}$$

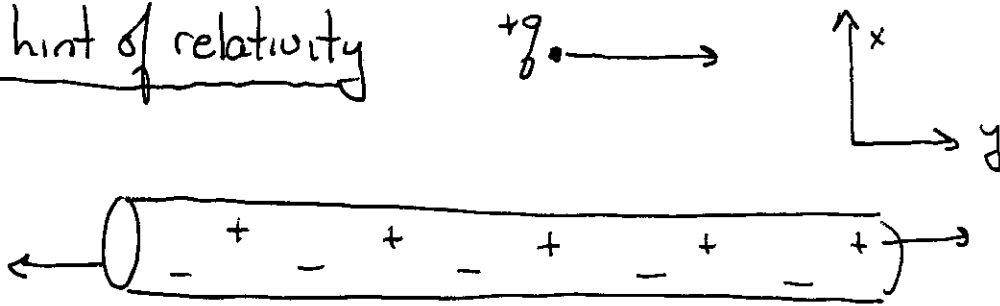
Magnetic force law $\vec{F} = q \frac{\vec{v}}{c} \times \vec{B}$



x \vec{B} out onto page

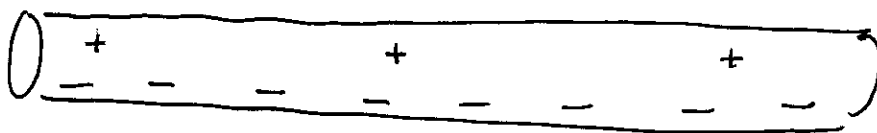
positive charge moving with current deflected toward current - like currents attract!

first hint of relativity



frame in which charge is moving
Lorentz force in -x direction

+q •



frame in which charge is at rest
electric force in -x direction

Magnetostatics \vec{J}, \vec{B} time-independent

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In general, charge is conserved

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0 \quad \text{continuity equation}$$

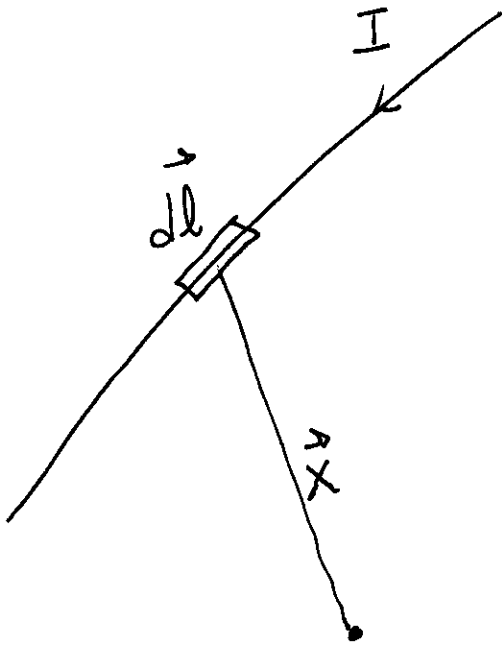
$$\int_V d^3x \left(\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} \right) = \frac{d}{dt} Q + \oint_S da \hat{n} \cdot \vec{J} = 0$$

change of charge in V flux out of V thru S

$$= 0$$



in magnetostatics, $\vec{\nabla} \cdot \vec{J} = 0$



$$d\vec{B} = \frac{I}{c} \frac{d\vec{l} \times \vec{r}}{|\vec{r}|^3}$$

but only makes sense for closed loop ($\vec{\nabla} \cdot \vec{J} = 0$)

$$[I] = \text{statcoulombs/s}$$

$$= \text{statamps}$$

$$[\vec{J}] = \text{statcoulombs/s/cm}^2$$

$$= \text{statamps/cm}^2$$

$$I d\vec{l} \rightarrow \vec{J}(\vec{x}') d^3x'$$

$$\vec{J} \sim I \delta^2(\vec{x}')$$

δ^2 transverse to wire

$$\vec{B}(\vec{x}) = \frac{1}{c} \int \vec{J}(\vec{x}') \times \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} d^3x'$$

$$= \frac{1}{c} \int \left(\vec{\nabla}'_{\vec{x}} \frac{1}{|\vec{x} - \vec{x}'|} \right) \times \vec{J}(\vec{x}') d^3x'$$

$$\vec{B}(\vec{x}) = \vec{\nabla} \times \left(\frac{1}{c} \int \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x' \right)$$

$$= \vec{\nabla} \times \vec{A} \quad \Rightarrow \quad \vec{\nabla} \cdot \vec{B} = 0 \text{ automatically}$$

$$\vec{A}(\vec{x}) = \frac{1}{c} \int \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x' + \vec{\nabla} \chi$$

gradient of scalar
gauge freedom

Consider

$$\begin{aligned} \vec{\nabla} \times \vec{B} &= \vec{\nabla} \times (\vec{\nabla} \times \vec{A}) \\ &= \vec{\nabla}(\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} \end{aligned}$$

$$\vec{\nabla}(\vec{\nabla} \cdot \vec{A}) = \frac{1}{c} \vec{\nabla} \left(\vec{\nabla} \cdot \int \frac{\vec{J}(\vec{x}') d^3 x'}{|\vec{x} - \vec{x}'|} \right)$$

$$= \frac{1}{c} \vec{\nabla} \int d^3 x' \vec{J}(\vec{x}') \cdot \vec{\nabla} \frac{1}{|\vec{x} - \vec{x}'|}$$

$$= -\frac{1}{c} \vec{\nabla} \int d^3 x' \vec{J}(\vec{x}') \cdot \vec{\nabla}' \frac{1}{|\vec{x} - \vec{x}'|}$$

$$= \frac{1}{c} \vec{\nabla} \int d^3 x' (\vec{\nabla}' \cdot \vec{J}(\vec{x}')) \frac{1}{|\vec{x} - \vec{x}'|} + \text{body terms}$$

$$= 0$$

second term: use $\nabla^2 \frac{1}{|\vec{x} - \vec{x}'|} = -4\pi \delta(\vec{x} - \vec{x}')$

$$\vec{\nabla} \times \vec{B} = -\frac{1}{c} \int d^3 x' \vec{J}(\vec{x}') \nabla^2 \frac{1}{|\vec{x} - \vec{x}'|}$$

$$= \frac{4\pi}{c} \vec{J}$$