

# Lagrangian of Relativistic Charged Particles in external EM fields

Lorentz force

$$\frac{du^\alpha}{d\tau} = \frac{q}{m} F^{\alpha\beta} u_\beta \quad u^\alpha = (\gamma_c, \vec{\gamma v})$$

$$u^\alpha = \frac{p^\alpha}{mc}$$

what about Lagrangian or  
Hamiltonian mechanics to deduce  
equations of motion?

note: my particle physics  
convention for  $p^\alpha$  ( $E, \vec{p}_c$ )

$$L[q_i(t), \dot{q}_i(t), t] \quad \begin{matrix} \text{generalized coordinates} \\ q, \dot{q} \end{matrix}$$

$$A = \int_{t_1}^{t_2} L dt \rightarrow \text{path is an extremum for } A \quad (\text{principle of least action})$$

$$\delta A = 0$$

leads to Euler-Lagrange equations of motion

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$

$$L = T - V \quad \begin{matrix} \text{non-relativistic} \\ \text{kinetic energy} \\ \text{potential} \end{matrix} \quad \text{only}$$

~~Max~~

want  $A$  to be Lorentz invariant scalar

$$\text{change to } A = \int_{\tau_1}^{\tau_2} \gamma L d\tau \quad \text{proper time}$$

$\therefore \gamma L$  must also be Lorentz invariant

$$p_i = \frac{\partial L}{\partial \dot{q}_i} \quad \text{generalized momentum}$$

when  $L$  does not depend on  $q_i$  then  
 $\frac{d}{d\tau}(p_i) = 0$  conservation of generalized  
momentum coordinate

for free particle  $p_i = \gamma_m v_i \Rightarrow \text{want } \frac{\partial L_{\text{free}}}{\partial v_i} = \gamma_m v_i$

$L_{\text{free}}$  is the T part

kinetic energy not suitable  
for relativistic

$$\text{if } L_{\text{free}} = -\frac{mc^2}{\gamma} = -mc^2 \sqrt{1-\beta^2} = -\frac{mc^2}{\gamma}$$

$$\text{then } \frac{\partial L_{\text{free}}}{\partial v_i} = -\frac{1}{2} mc^2 \left( \frac{\gamma}{\sqrt{1-\beta^2}} \right) \left( -\frac{2v}{c^2} \right) = \gamma m v$$

and clear that  $\gamma L = -mc^2$  is Lorentz invariant

now add the V part for EM

$V = e\Phi$  in non-relativistic limit

$$L_{\text{int}} \rightarrow L_{\text{int}}^{NR} = -e\Phi$$

for relativistic and covariant formulation:

$$L_{\text{int}} = -\frac{e}{\gamma} U_\alpha A^\alpha$$

~~A<sup>α</sup> (Φ, A)~~

$$A^\alpha = \left( \frac{\Phi}{c}, \vec{A} \right)$$

$$L_{\text{int}} = -e\Phi + e\vec{V} \cdot \vec{A}$$

$$U_\alpha = (\gamma_c, -\gamma\vec{v})$$

full Lagrangian relativistic charged particle in EM field is:

$$L = -\frac{mc^2}{\gamma} - e\Phi + e\vec{V} \cdot \vec{A}$$

does this generate the correct  
equations of motion?

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$

~~$\frac{\partial L}{\partial t} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$~~

$$\frac{\partial L}{\partial \dot{q}_i} = e \vec{\nabla} \left[ \vec{v} \cdot \vec{A} - \phi \right] = e \left[ (\vec{v} \cdot \vec{\nabla}) \vec{A} + \vec{v} \times (\vec{\nabla} \times \vec{A}) \right] - e \vec{\nabla} \phi$$

$$\vec{\nabla}(\vec{a} \cdot \vec{b}) = (\vec{a} \cdot \vec{\nabla}) \vec{b} + (\vec{b} \cdot \vec{\nabla}) \vec{a} + \vec{a} \times (\vec{\nabla} \times \vec{b}) + \vec{b} \times (\vec{\nabla} \times \vec{a})$$

and  $\vec{v}$  is just ~~not~~ a vector, does not depend on  $q_i$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \vec{v}} \right) = \frac{d}{dt} \left[ \gamma m \vec{v} + \frac{\partial}{\partial \vec{v}} (e \vec{v} \cdot \vec{A}) \right] = \frac{d}{dt} (\gamma m \vec{v}) + e \frac{d \vec{A}}{dt}$$

note:  $\frac{d \vec{A}}{dt} = \frac{\partial \vec{A}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{A}$  "convective derivative"

$$\therefore \frac{d}{dt} (\gamma m \vec{v}) = e \left[ -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t} + \vec{v} \times \underbrace{(\vec{\nabla} \times \vec{A})}_{\vec{B}} \right]$$

Lagrangian for charged particle in EM field has canonical momentum or conjugate momentum or generalized to position  $\vec{x}$

$$\overset{\cancel{P}}{P}_i = \frac{\partial L}{\partial v_i} = \gamma m v_i + e A_i = p_i + e A_i$$

$$\boxed{\vec{P}_i = \vec{p} + e \vec{A}}$$